

Quantum random number generator based on the photon number decision of weak laser pulses

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We propose an approach to realize a quantum random number generator (QRNG) based on the photon number decision of weak laser pulses. This type of QRNG can generate true random numbers at a high speed and can be adjusted to zero bias conveniently, thus is suitable for the applications in quantum cryptography.

Random numbers are essential in a very wide application range, such as statistical sampling [1], computer simulations [2], randomized algorithm [3] and cryptography[4]. In the application of quantum cryptography, true random numbers are required for the secure key distribution. Current theory implies that the only way to realize a random number generator (RNG) which can be scientifically proved to be undeterministic is to use the intrinsic randomness of quantum decisions, for the occurrence of each possible result is unpredictable. Some practical methods to realize a quantum random number generator (QRNG) have been proposed: single photons incident on a 50:50 beam splitter [5]; polarized single photons incident on a rotatable polarizing beam splitter [5] or Fresnel multiple prism [6]; and utilizing the random time intervals between photon emissions of semiconductors [7]. Since no *practical* single photon source exists nowadays, the QRNGs based on single photons usually use weak laser source to approximate the single photon source. Consequently, the generation rate of random numbers is restricted by the probability of single-photon component in laser pulse. In this Letter, we propose an approach for randomness generation based on the photon number decision of weak laser pulses. This type of QRNG can generate one random bit for each random event and it can be conveniently adjusted to the state of generating ones and zeros with equal probabilities. Besides, it has a more compact set up. Those advantages make it suitable for the applications of quantum cryptography.

We define a random bit generator as a device which produces bits independently of each other and with equal probabilities of ones and zeros, i.e., $p(0) = p(1) = 0.5$. Normally, the photon number distribution of weak laser pulses is Poissonian [8]. Since the photon number distribution of partially absorbed light follows a *Bernoulli transform* of the initial field [9], the detected photon number distribution of weak laser pulses follows

$$P_\eta(n) = \frac{(\eta\lambda)^n e^{-\eta\lambda}}{n!}, \quad (1)$$

where λ is the mean photon number of the weak laser pulses, η is the detection efficiency of the single photon detector.

Experimentally, we use an avalanche photodiode (APD) operating in gated mode in the measurement, which does not distinguish the photon numbers above zero photons. In this situation, we get the result ‘0’ when the pulse contains no photon, and the result ‘1’ when above zero photon. Hence, the probabilities of getting results of ‘0’ and ‘1’ are $P_\eta(0) = e^{-\eta\lambda}$, $P_\eta(1) = 1 - e^{-\eta\lambda}$, respectively. We then have $P_\eta(0) = P_\eta(1) = 0.5$, when $\eta\lambda = 0.693$. Since η is a specification of the detector, we can simply adjust λ to set $\eta\lambda$ to the proper value (0.693). Concerning that the probabilities of generating ones and zeros are equal and each generation is independent, the outcome of the QRNG is true random.

The experimental set up of our QRNG is shown in Fig. 1. We use a pulsed laser source (PLS, id300, produced by id Quantique) to generate laser pulses of 300 ps at 1550 nm according to external trigger. First, the controlling system generates an NIM signal of 1 MHz to trigger id300. The emerging laser pulses are coupled into single mode fiber (SMF) and then pass through a mechanically adjustable attenuator. Finally, the pulses are detected by the single photon detector module (SPDM, id200, produced by id Quantique). The module of id200 is based on an InGaAs APD working in gated mode, where a voltage pulse is applied to raise the bias above breakdown upon triggering. If there are photons detected during a gate, the SPDM will output a logic ‘1’ signal after the gate, otherwise the response will be logic ‘0’. We set the dead time of id200 as zero and the gate width as 2.5 ns. The controlling system generates a TTL trigger signal for id200 with a proper delay from the trigger of id300. The dark count rate is measured to be 10^{-5} in experiment, and the detection efficiency is no less than 10 percent according to the features provided by id Quantique. The average power of id300 at 1 MHz is -35 ± 1 dBm according to the specifications. Taking the detection efficiency of id200 as 0.1, the average photon number after the attenuation should be 6.93. Since the transmittance of the attenuator can be continuously adjusted from 0 to -30 dB, the probabilities of generating ones and zeros can be practically adjusted to be equal. The output of the SPDM is recorded and transferred to PCI-7300A (PC interfaced data acquisition board, produced by ADLINK Technology Inc.) by the controlling system. In order to eliminate the errors due to the clock drift between PCI-7300A and id200, the controlling system accompanies the data signal with a synchronizing clock. We develop the controlling system based on a chip of

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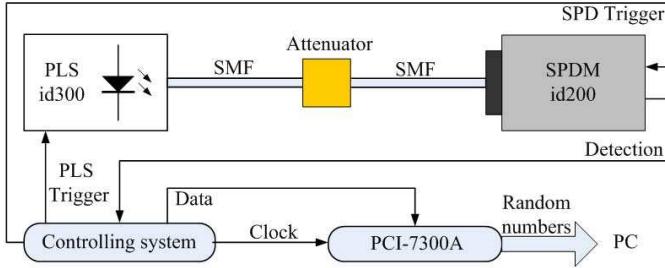


FIG. 1: (color online) Schematic experiment set up for our QRNG.

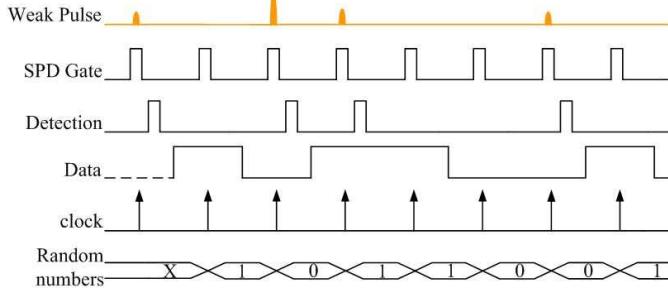


FIG. 2: (color online) The operation (clock to clock) of our QRNG.

Cyclone II EP2C5T144C8 (FPGA, produced by Altera). The peak power of the laser source is 0 dBm according to the specifications provided by id Quantique. It is much less than the 10 dBm damage level of the SPDM. First, we apply a relatively weak attenuation to the emitted laser pulses so that almost every pulse can be detected by the SPDM. We scan the delay between the trigger of the laser and the detection gate to find the maximum count rate, which means the detection gate catch the laser pulses. After finishing the synchronization of the experimental system, we adjust the mechanical attenuator delicately to make the photon count rate approach to half of the repetition rate of the laser pulses. We fix the attenuator at the best point and find that our QRNG is thus set up. The operation (clock to clock) of our QRNG is illustrated schematically in Fig. 2. The QRNG described above can be compared with an ideal random number generator, which produces ones and zeros with equal probability and every bit is totally independent of the previous ones. According to [10], the serial correlation coefficient a_k of the sequence $Y_1 \dots Y_N$ with lag $k \geq 1$ is defined as

$$a_k = \frac{\sum_{i=1}^{N-k} (Y_i - \bar{Y})(Y_{i+k} - \bar{Y})}{\sum_{i=1}^N (Y_i - \bar{Y})^2}, \quad (2)$$

where N and \bar{Y} are the length and the mean of the sequence, respectively. The coherence time of the laser is about 13ps according to its spectral width of 0.6nm. Time interval between two consecutive laser pulses is much larger than the coherence time for the repetition rate within 10 GHz. Thus the photon decision of each pulse is totally independent. When the probabilities of generating ones and zeros are equal, a_k should approach to zero. However, there always exists the after pulse effect for real SPDs, which means that the SPD

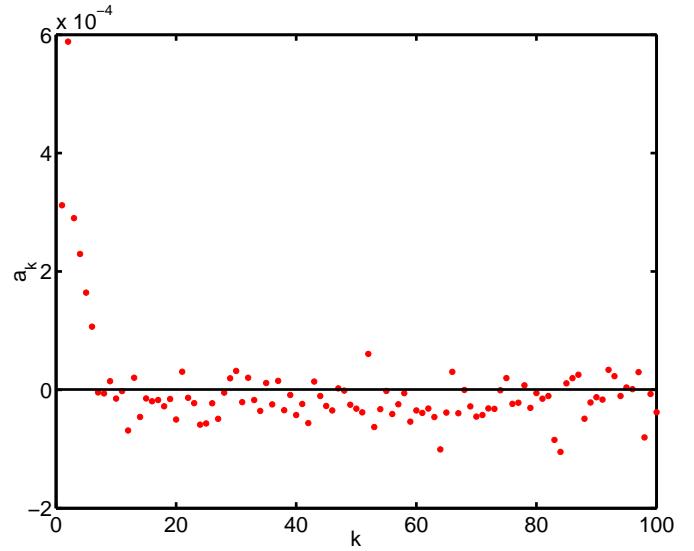


FIG. 3: (color online) Serial correlation coefficient (a_k) of a typical sequence of 10^9 bits, with the lag k ranging from 1 to 100.

would generate an electrical pulse with a certain probability within a gate period if it detects photons in the previous gate. The after pulse effect obviously introduces correlations to the binary sequence generated by the SPD. For a typical sequence of 10^9 bits, we calculate a_k with k ranging from 1 to 100, the results of which are plotted in Fig. 3. It can be seen that a_k is relatively high for $k \leq 6$ due to the after pulse effect. Hence, for the true randomness, we should select the random bit with an interval of more than 6 bits. Another modification of our QRNG is the non-equal probabilities of ones and zeros. It can be described by the bias b which is defined as $b \equiv P(0) - 1/2$. The bias of our QRNG varies with the detected mean photon number, which is dominated by the detection efficiency of SPD, the transmittance of attenuator, and the power of laser. The average of detected photon number λ_d can be written as

$$\lambda_d = \frac{P_{avg} T \eta}{h \nu_0 f_{rep}}, \quad (3)$$

where P_{avg} is the average optical power, ν_0 is the central frequency of the pulsed laser, h is Planck's constant, T is the transmittance of the attenuator, η is the detection efficiency of the SPD, and f_{rep} is the repetition rate of the laser. In experiments, a feed forward loop can be built to ensure the long term stability of λ_d . Thus, the bias is restricted to zero physically. Further more, the original bits could be unbiased with appropriate mathematical procedures [11, 12].

So far, there is no generally accepted definition of absolute randomness. In the applications of quantum cryptography, the most desired feature of RNG is its impossibility of description or prediction. This can only be proved by recording the outcome of the RNG for infinite time. However, only finite samples are practically possible. Many empirical methods are thus proposed to test RNG with sequences of finite length. Though these empirical methods are not sufficient for the test of the true randomness, they may detect some imperfections

TABLE I: Results of ENT for a typical sequence of 10^9 bits

Entropy = 1.000000 bits per bit.
Optimum compression would reduce the size of this 1025999992 bit file by 0 percent.
Chi square distribution for 1025999992 samples is 2.07, and randomly would exceed this value 15.00 percent of the times.
Arithmetic mean value of data bits is 0.5000 (0.5 = random).
Monte Carlo value for P_i is 3.140925340 (error 0.02 percent).
Serial correlation coefficient is 0.000312 (totally uncorrelated = 0.0).

TABLE II: Results of DIEHARD for a typical sequence of 10^9 bits

Birthday Spacings	0.443957
Overlapping Permutations	0.467282
Ranks of 31×31 Matrices	0.988764
Ranks of 32×32 Matrices	0.364029
Ranks of 6×8 Matrices	0.359226
Monkey Tests on 20-bit Words	0.31820
Monkey Test OPSO	0.2027
Monkey Test OQSO	0.5306
Monkey Test DNA	0.3234
Count 1's in Stream of Bytes	0.543339
Count 1's in Specific Bytes	0.684855
Parking Lot Test	0.165163
Minimum Distance Test	0.501530
Random Spheres Test	0.356525
The SQUEEZE Test	0.716755
Overlapping Sums Test	0.437118
Runs Test (up)	0.777892
Runs Test (down)	0.854107
The Craps Test No. of wins	0.808609
The Craps Test throws/game	0.511087

of the random numbers. We choose two batteries of statistical tests to evaluate our QRNG: ENT and DIEHARD, which we consider to be sufficient in qualifying the device for its use in the experiment.

ENT [13] is a series of basic statistical tests which evaluate the random sequence in some elementary features such as

the equal probabilities of ones and zeros and the serial correlation. The testing results of a typical sequence of 10^9 bits are presented in Table I. From the results, we can see that our QRNG generates ones and zeros with almost equal probability. The serial auto-correlation coefficient is of the order of 10^{-4} , which is due to the after pulse effect. In the standard ENT test, the Monte Carlo estimation for π actually evaluates the uniform distribution of blocks of 48 bits.

To further exploit some subtle imperfections hidden in our QRNG, we test the sample sequence using DIEHARD [14]. DIEHARD is widely considered as one of the best strengthened randomness testing battery because it is most sensitive to various problems possible in pseudo RNG. It consists of 15 tests with outcome of one or more p -values. According to the instruction of the testing suit, a sequence could not be considered as random if p -value is less than 0.01 or greater than 0.99 for six or more places. The testing results of the sequence of 1×10^9 bits are shown in Table II, from which one can find that our QRNG generates true random numbers.

We present an approach of QRNG based on the photon number decision of weak laser pulses. The realization of it consists of a pulsed laser source, a flexible attenuator, a single photon detector, and some circuits used for controlling and data acquisition. This type of QRNG has advantages for the application in quantum cryptography. It can generate random numbers at a high speed that is limited only by the recovery time of the single photon detector. If the device is realized with fast single photon detectors, e.g., the ones based on silicon APD, the generation rate of random numbers is hopefully increased to GHz or higher. In addition, this type of QRNG can be more compact for it needs only one APD photon counter.

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